

Fourier Based MLFMM for the Helmholtz Equation

Cris Cecka Eric Darve

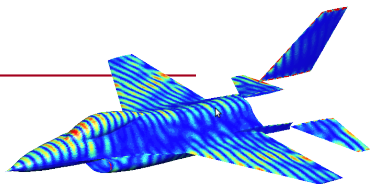
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Institute for Computational and Mathematical Engineering
Stanford University

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Solving Helmholtz



The Helmholtz equation over $\Omega \subset \mathbb{R}^3$,

$$(\nabla^2 + \kappa^2)u(\mathbf{x}) = \rho(\mathbf{x})$$

has solution

$$u(\mathbf{x}) = \int_{\Omega} \frac{e^{i\kappa|\mathbf{x}-\mathbf{y}|}}{4\pi|\mathbf{x}-\mathbf{y}|} \rho(\mathbf{y}) d\mathbf{y}$$

Taking a set of point sources $\{\mathbf{y}_j\}_{j=1}^N$ and field points $\{\mathbf{x}_i\}_{i=1}^N$,

$$u(\mathbf{x}_i) = \sum_{\mathbf{y}_j \neq \mathbf{x}_i} \frac{e^{i\kappa|\mathbf{x}_i-\mathbf{y}_j|}}{4\pi|\mathbf{x}_i-\mathbf{y}_j|} \rho_j$$

requiring $\mathcal{O}(N^2)$ operations. MLFMM reduces this to $\mathcal{O}(N \log N)$.



Low Rank Approximation

The Helmholtz MLFMM is based on the plane wave expansion

$$\frac{e^{i\kappa|\mathbf{r}_0+\mathbf{r}|}}{|\mathbf{r}_0+\mathbf{r}|} = \lim_{\ell \rightarrow \infty} \int_{S^2} e^{i\kappa\langle \hat{\mathbf{s}}, \mathbf{r} \rangle} T_{\ell, \mathbf{r}_0}(\hat{\mathbf{s}}) dS(\hat{\mathbf{s}})$$

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y_j •



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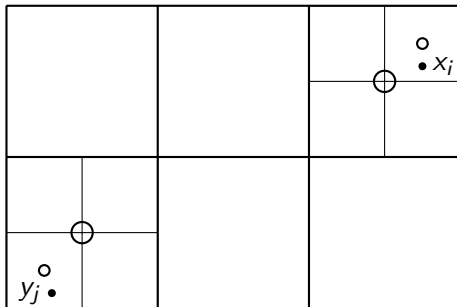


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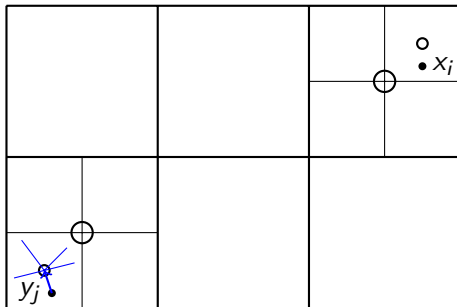


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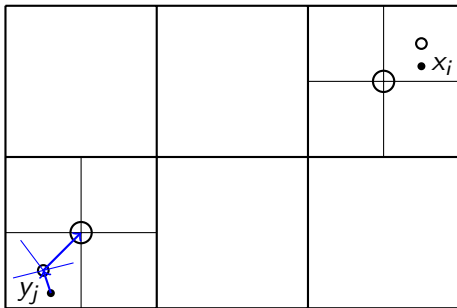


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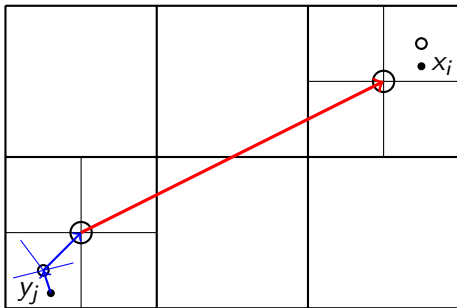


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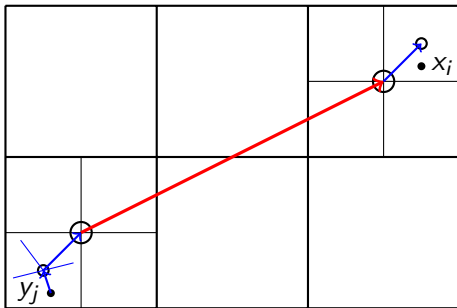


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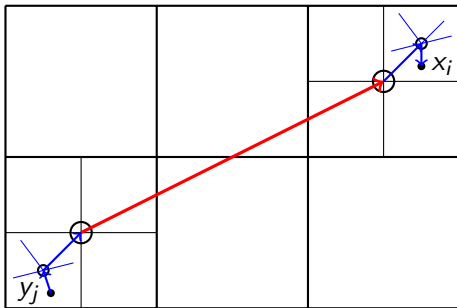


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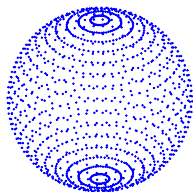
Typical Helmholtz MLFMM

- Field representation on S^2 - Spherical quadrature



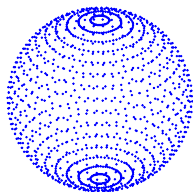
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 - Exactly integrates all Y_n^m , $|m| \leq n$, $n < N$.



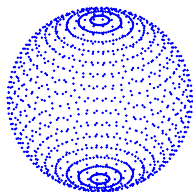
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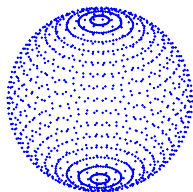
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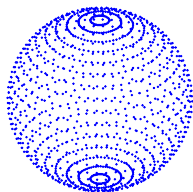
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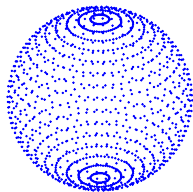
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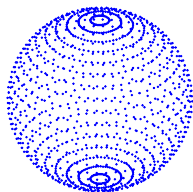
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- Moral: Want better interpolation/anterpolation.



Fourier Based Helmholtz MLFMM

Instead, use a 2D Fourier basis:

$$\int_{S^2} e^{i\kappa\langle\hat{\mathbf{s}},\mathbf{r}\rangle} T_{\ell,\mathbf{r}_0}(\hat{\mathbf{s}}) dS(\hat{\mathbf{s}}) = \int_0^{2\pi} \int_0^{2\pi} e^{i\kappa\langle\hat{\mathbf{s}},\mathbf{r}\rangle} T_{\ell,\mathbf{r}_0}^{\mathbf{s}}(\hat{\mathbf{s}}) d\theta d\phi$$

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- Uniform quadrature in $[0, 2\pi] \times [0, \pi]$ (ϕ -symmetry)



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 - False.
 - Can be computed efficiently.
 - Can retain diagonality.



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Consider a 1D analogue:

$$\mathcal{I} = \int_0^{2\pi} \left[\sum_{n=-N}^N a_n e^{in\phi} \right] \left[\sum_{m=-M}^M b_m e^{im\phi} \right] d\phi \quad M \gg N$$



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- Remove frequencies of $T_{\ell, \mathbf{r}_0}^{\mathbf{s}}(\hat{\mathbf{s}})$ which do not contribute.
 - Make $T_{\ell, \mathbf{r}_0}^{\mathbf{s}}(\hat{\mathbf{s}})$ band-limited with $\mathcal{O}(\kappa |r|)$ frequencies.
 - Easy since T_{ℓ, \mathbf{r}_0} has bandwidth ℓ and

$$\mathcal{F}(m; |\sin(\theta)|) = \frac{(-1)^m + 1}{\pi(1 - m^2)}$$



Improved Error Analysis

With exact integration and interpolation, only two sources of error:

$$\begin{aligned} \frac{e^{i\kappa|\mathbf{r}_0+\mathbf{r}|}}{|\mathbf{r}_0+\mathbf{r}|} &\stackrel{1}{\rightarrow} \iint_0^{2\pi} e^{i\kappa\langle\hat{\mathbf{s}},\mathbf{r}\rangle} T_{\ell,\mathbf{r}_0}^{\mathbf{s}}(\hat{\mathbf{s}}) d\theta d\phi \\ &\stackrel{2}{\rightarrow} \sum_k \omega_k e^{i\kappa\langle\hat{\mathbf{s}}_k,\mathbf{r}\rangle} T_{\ell,\mathbf{r}_0}^{\mathbf{s},L}(\hat{\mathbf{s}}_k) \end{aligned}$$

- Error 1 - Gegenbauer series truncation.
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 - Account for error in $T_{\ell,\mathbf{r}_0}^{\mathbf{s}} \rightarrow T_{\ell,\mathbf{r}_0}^{\mathbf{s},L}$.
 - Account for error in sampling $e^{i\kappa\langle\hat{\mathbf{s}}_k,\mathbf{r}\rangle}$ in finite quadrature.



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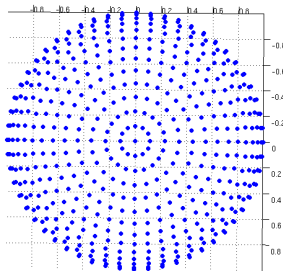
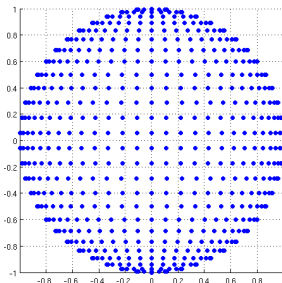
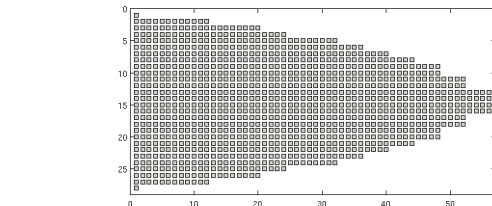
$$\begin{aligned} \frac{e^{i\kappa|\mathbf{r}_0+\mathbf{r}|}}{|\mathbf{r}_0+\mathbf{r}|} &\xrightarrow{1} \iint_0^{2\pi} e^{i\kappa\langle\hat{\mathbf{s}},\mathbf{r}\rangle} T_{\ell,\mathbf{r}_0}^{\mathbf{s}}(\hat{\mathbf{s}}) d\theta d\phi \\ &\xrightarrow{2} \sum_k \omega_k e^{i\kappa\langle\hat{\mathbf{s}}_k,\mathbf{r}\rangle} T_{\ell,\mathbf{r}_0}^{\mathbf{s},L}(\hat{\mathbf{s}}_k) \end{aligned}$$

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 - Want to find the smallest quadrature for a given accuracy.
 - Spectral analysis yields constructive algorithms for quadrature.



Quadrature Example

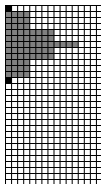
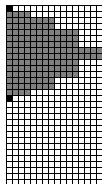
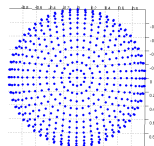
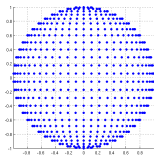
With $\kappa |\mathbf{r}| = 10$, a 1042 point quadrature is created:



Interpolation/Anterpolation Between Quadratures

Interpolate/Anterpolate from quadrature Q to Q' . Let

$$\mathcal{N}_\phi = \max\left(\max_{1 \leq j \leq N_\theta} N_\phi(\theta_j), \max_{1 \leq j \leq N'_\theta} N'_\phi(\theta_j) \right)$$

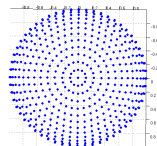
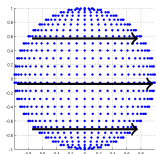
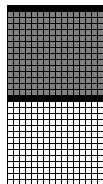
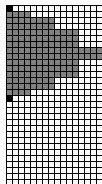


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- 1 Interpolate $N_\phi(\theta_j)$ to \mathcal{N}_ϕ .

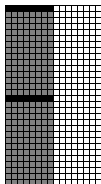
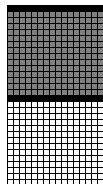
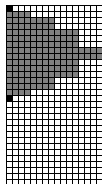
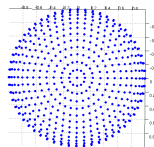
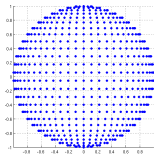


Interpolation/Anterpolation Between Quadratures

Interpolate/Anterpolate from quadrature Q to Q' . Let

$$\mathcal{N}_\phi = \max\left(\max_{1 \leq j \leq N_\theta} N_\phi(\theta_j), \max_{1 \leq j \leq N'_\theta} N'_\phi(\theta_j) \right)$$

- 1 Interpolate $N_\phi(\theta_j)$ to \mathcal{N}_ϕ .
- 2 Wrap the data to construct the θ -periodic data.

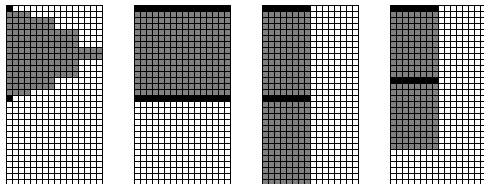
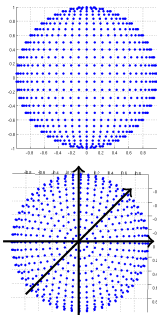


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- 1 Interpolate $N_\phi(\theta_j)$ to \mathcal{N}_ϕ .
- 2 Wrap the data to construct the θ -periodic data.
- 3 Interpolate from length N_θ to N'_θ .

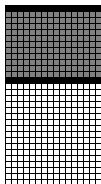
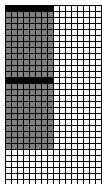
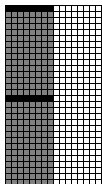
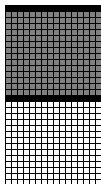
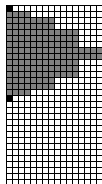
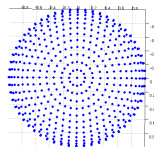
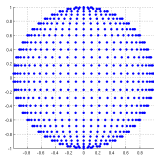


Interpolation/Anterpolation Between Quadratures

Interpolate/Anterpolate from quadrature Q to Q' . Let

$$\mathcal{N}_\phi = \max\left(\max_{1 \leq j \leq N_\theta} N_\phi(\theta_j), \max_{1 \leq j \leq N'_\theta} N'_\phi(\theta_j) \right)$$

- 1 Interpolate $N_\phi(\theta_j)$ to \mathcal{N}_ϕ .
- 2 Wrap the data to construct the θ -periodic data.
- 3 Interpolate from length N_θ to N'_θ .
- 4 Unwrap the data to construct the ϕ -periodic data.

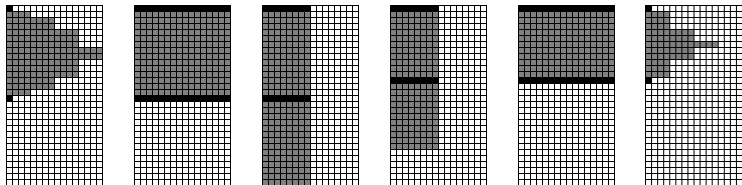
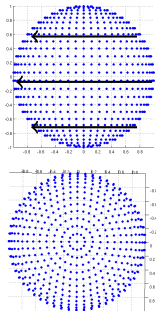


Interpolation/Anterpolation Between Quadratures

Interpolate/Anterpolate from quadrature Q to Q' . Let

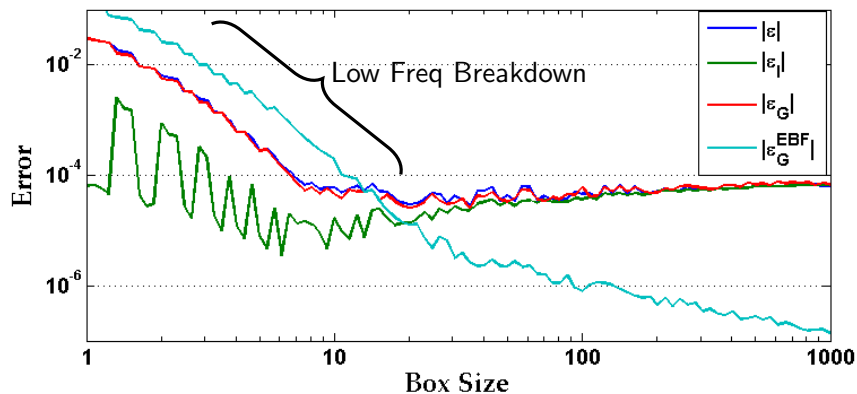
$$\mathcal{N}_\phi = \max\left(\max_{1 \leq j \leq N_\theta} N_\phi(\theta_j), \max_{1 \leq j \leq N'_\theta} N'_\phi(\theta_j)\right)$$

- 1 Interpolate $N_\phi(\theta_j)$ to \mathcal{N}_ϕ .
- 2 Wrap the data to construct the θ -periodic data.
- 3 Interpolate from length N_θ to N'_θ .
- 4 Unwrap the data to construct the ϕ -periodic data.
- 5 Anterpolate from length \mathcal{N}_ϕ to $N'_\phi(\theta_j)$.

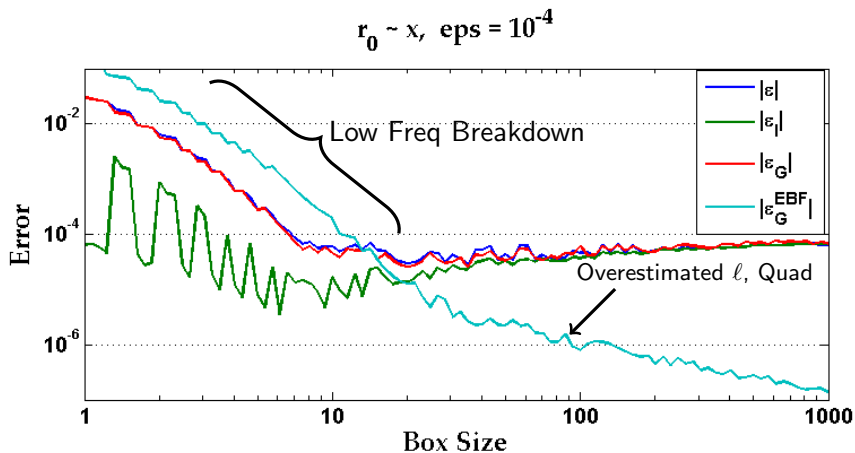


Error Results

$$r_0 \sim x, \text{ eps} = 10^{-4}$$

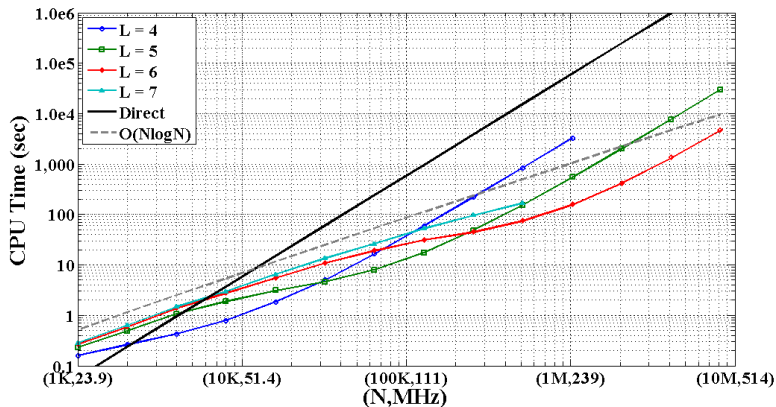


Error Results



Asymptotic Running Time

With accuracy approximately 10^{-6} ,



Moral of the Story

Reasons to use the Fourier Based FMM:

- Exact interpolation/interpolation steps
- Fast with widely available, highly optimized FFT libraries
- Easier to implement - (Uniform quadrature, FFTs, etc)
- All error can be accounted for in the precomputation stage
 - Strict error bounds to be computed before the computation.
 - Error bounds are used to determine truncation and quadrature.

